# Algorithm of evaluation of the unstable equilibrium state of dynamically tuned gravimeter 

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#### Abstract

The article describes the development and investigation of the algorithm of valuation of the dynamically tuned gravimeter state using the least-squares method on the basis of information about the movement of sensors regarding the unstable equilibrium; evaluation errors caused by the inadequacy of the accepted approximation of the model and the real signal, and the errors caused by the kinematic nonlinearity have been also studied and described.


Index Terms- gravimetric aviation system, dynamically tuned gravimeter, output signal, evaluation errors, algorithm.

## 1 Introduction

he information about the Earth's gravitational field anomalies is used in inertial navigation and in the search of mineral resources. Those gravimeters (among the known ones) [1-2] that are used in aviation gravimetric systems have the lowest accuracy, since they are the most disadvantaged in terms of measuring and eliminating external factors that influence the measuring signal of the gravimeter. Therefore, increasing the accuracy and speed of the dynamically tuned gravimeter of the gravimetric aviation system (GAS) with the automatic processing of information is conditioned by the need of establishing the effective and easy-to-implement algorithms of valuation of the dynamically tuned gravimeter (DTG) state.

In the analyzed literature [3-8] the modifications of the optimal and suboptimal algorithms of discrete signal filtering are investigated, but the studies of the DTG accuracy [3-8] do not take into account the impact of gyrogravimeter errors caused by nonlinear distortions of gyroscope motion path, by inequality of index of precession damping vibrations to zero due to the influence of viscous friction, by nonisochism of precession fluctuations, by the mismatch of precession vibrations angular frequency used in the evaluation algorithms to the gyroscope precession fluctuations frequency, by the barriers that distort the gyroscope motion law. The impact of these errors, if they are not taken into account, significantly influences the output signal of DTG and reduces the measurement accuracy of Earth's gravitational field anomalies.

Since the task of optimal filtration is the task of the state evaluation in stochastic terms, we will use the term "state evaluation" in the future. Taking into account the interconnection of the methods of least squares (MLS) and optimal Kalman filter (KF), we will consider that MLS is the simplest approach to the state evaluation, and KF is the only optimal evaluation algorithm. Optimal Kalman filter was investigated

[^0]in [9]. The rest of the evaluation algorithms are intermediate between MLS and KF, so it is reasonable to further investigate only these two algorithms.

The aim of this paper is to develop DTG state evaluation algorithm under the conditions of orientation of its sensitivity axis to the north and to the south by the method of least squares. The achievement of this goal will allow the improvement of the accuracy and speed of DTG GAS measurement by eliminating previously uncompensated errors.

To achieve this goal it is necessary to solve the following problems:

1) to develop DTG state evaluation algorithm under the conditions of orientation of its sensitivity axis to the north and to the south (stable and unstable equilibrium) on the basis of the method of least squares;
2) to investigate evaluation errors, to determine their dependence on the perturbation parameters, the initial conditions of dynamically tuned gravimeter motion state, sensitive element (SE);
3) to analyze the evaluation errors caused by the evaluation algorithms sensitivity to the discrepancies of the accepted model and real output signal of DTG;
4) to investigate the influence of distortion of DTG motion law seen as the typical distortions and measurement channel angle sensor (AS) noise.

To solve this problem, DTG $[10,11]$ was applied . Output signal of dynamically tuned gravimeter, which is taken from the AS is fed into the computer. The computer determines the value of the errors using the algorithm described in this article and provides automatic compensation of these errors in the output signal of the gyrogravimeter. Then the adjusted signal of the gyrogravimeter along with the signals from other subsystems of GAS is used to calculate $\Delta \mathrm{g}$.

Main part. As in [11], sensing element motion observed by the angle sensor is represented by the function

$$
\begin{equation*}
\alpha(t)=R^{S}+\alpha_{1}^{*}(t)+\varepsilon(t) \tag{1}
\end{equation*}
$$

where $R^{S}$ is a true angle between AS zero and the direction to the south; $\alpha_{1}^{*}(t)$ is current angular position of SE relative to

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the direction to the south, which is determined as the solution of equation (1) in the case of replacement $\alpha_{1}(t)=\pi+\alpha_{1}^{*}(t)$. For this case the differential equation (1) can be written as

$$
\begin{equation*}
\ddot{\alpha}_{1}^{*}(t)+2 \xi_{1} \dot{\alpha}_{1}^{*}-\omega_{0}^{2} \sin \alpha_{1}^{*}(t)=0 \tag{2}
\end{equation*}
$$

We will omit index "*" for the ease of writing. When $\sin \alpha \approx \alpha$, the solution of equation (2) is described by the expression

$$
\begin{equation*}
\alpha_{1}(t)=e^{-\xi_{1} t}\left(c_{10} \operatorname{sh} \omega_{0} t+c_{20} \operatorname{ch} \omega_{0} t\right) \tag{3}
\end{equation*}
$$

where $c_{10}, c_{20}$ are determined by the initial conditions at $t=0$. Based on the expressions (1) and (2), when $\xi_{1}=0$, the SE motion model looks like

$$
\begin{equation*}
\alpha(t)=\hat{R}^{s}+\hat{c}_{1} \operatorname{sh} \omega_{0} t+\hat{c}_{2} \operatorname{ch} \omega_{0} t \tag{4}
\end{equation*}
$$

where $\hat{R}^{s}$ is the value of the angle between AS zero and the state of unstable equilibrium; $\hat{C}_{1}, \hat{c}_{2}$ is the values of $C_{10}, c_{20}$ respectively.

To develop the evaluation algorithm we generate the quadratic functional

$$
\begin{equation*}
F_{S}=\sum_{i=1}^{n}\left(\hat{R}^{S}+\hat{c}_{1} \operatorname{sh} \omega_{0} t_{i}+\hat{c}_{2} \operatorname{ch} \omega_{0} t_{i}-\alpha_{i}\right)^{2} \tag{5}
\end{equation*}
$$

The system of algebraic equations

$$
\begin{cases}\hat{R}^{s} \sum_{i=1}^{n} 1+\hat{c}_{1} \sum_{i=1}^{n} s h \omega_{0} t_{i} & +\hat{c}_{2} \sum_{i=1}^{n} \operatorname{ch} \omega_{0} t_{i}=\sum_{i=1}^{n} \alpha_{i}  \tag{6}\\ \hat{R}^{s} \sum_{i=1}^{n} \operatorname{sh} \omega_{0} t_{i}+\hat{c}_{1} \sum_{i=1}^{n} \operatorname{sh}^{2} \omega_{0} t_{i} & +\hat{c}_{2} \sum_{i=1}^{n} \operatorname{sh} \omega_{0} t_{i}=\sum_{i=1}^{n} \alpha_{i} \operatorname{sh} \omega_{0} t_{i} \\ \hat{R}^{s} \sum_{i=1}^{n} \operatorname{ch} \omega_{0} t_{i}+\hat{c}_{1} \sum_{i=1}^{n} \operatorname{sh} \omega_{0} t_{i} \operatorname{ch} \omega_{0} t_{i}+\hat{c}_{2} \sum_{i=1}^{n} \operatorname{ch}^{2} \omega_{0} t_{i}=\sum_{i=1}^{n} \alpha_{i} \operatorname{ch} \omega_{0} t_{i}\end{cases}
$$

meets the terms of the minimum functional.
System observation analysis (6), conducted the same way as in [9], shows that observation (solvability) of the system under study is provided on condition that

$$
\begin{equation*}
\lambda=\omega_{0} \Delta t \neq 0 ; \quad n \geq 3 \tag{7}
\end{equation*}
$$

After replacing of hyperbolic functions on the exponential, system (6) can be represented by a matrix equation

$$
\begin{equation*}
c_{\lambda, n}^{S} \hat{x}_{S}=z_{\lambda, n}^{S} \tag{8}
\end{equation*}
$$

where

$$
\begin{aligned}
& c_{\lambda, n}^{S}=\left|\begin{array}{ccc}
n & \frac{1-\mathrm{e}^{\lambda \mathrm{n}}}{1-e^{\lambda}} & \mathrm{e}^{\lambda(1-n)} \frac{1-\mathrm{e}^{\lambda \mathrm{n}}}{1-\mathrm{e}^{\lambda}} \\
\frac{1-e^{\lambda n}}{1-e^{\lambda}} & \frac{1-\mathrm{e}^{2 \lambda \mathrm{n}}}{1-e^{2 \lambda}} & \mathrm{n} \\
e^{\lambda(1-n)} \frac{1-e^{\lambda n}}{1-e^{\lambda}} & n & e^{2 \lambda(1-n)} \frac{1-\mathrm{e}^{2 \lambda \mathrm{n}}}{1-\mathrm{e}^{2 \lambda}}
\end{array}\right|, \\
& \hat{x}_{S}=\left[\hat{R}^{S} \hat{C}_{1} \hat{c}_{2}\right]^{T}, \quad z_{\lambda, n}^{S}=\left[S_{1}^{S} S_{2}^{S} S_{3}^{S}\right]^{T}=\left[\sum_{i=1}^{n} \alpha_{i} \sum_{i=1}^{n} \alpha_{i} e^{\lambda(i-1)} \sum_{i=1}^{n} \alpha_{i} e^{-\lambda(i-1)}\right]^{T} .
\end{aligned}
$$

Solving the system (8) by Kramer method, we obtain an expression for the state vector elements that are measured in the form of

$$
\begin{aligned}
& \hat{R}^{S}=\frac{A_{11}^{S} S_{1}^{S}+A_{22}^{S} S_{2}^{S}+A_{31}^{S} S_{3}^{S}}{\operatorname{det} c_{\lambda, n}^{S}} \\
& \hat{c}_{1}=\frac{A_{12}^{S} S_{1}^{S}+A_{22}^{S} S_{2}^{S}+A_{31}^{S} S_{3}^{S}}{\operatorname{det} c_{\lambda, n}^{S}} \\
& \hat{c}_{2}=\frac{A_{13}^{S} S_{1}^{S}+A_{23}^{S} S_{2}^{S}+A_{33}^{S} S_{3}^{S}}{\operatorname{det} c_{\lambda, n}^{S}},
\end{aligned}
$$

where $A_{i j}$ are algebraic complements, element $c_{i j}^{s}(i, j=\overline{1,3})$ of matrix $c_{\lambda, n}^{S} ; \quad \operatorname{det} c_{\lambda, n}^{S}$ is the determinant of matrix. Revealing more detailed the expression $\hat{R}^{s}$, we obtain

$$
\begin{align*}
& \hat{R}_{S}=\left\{S_{1}^{S}\left[\frac{1+e^{\lambda n}}{1+e^{\lambda}}+n e^{\lambda(n-1)} \frac{1-e^{\lambda}}{1-e^{\lambda n}}\right]-S_{2}^{S}-S_{3}^{S} e^{\lambda(n-1)}\right\} \times \\
& \times\left\{n\left[\frac{1+e^{\lambda n}}{1+e^{\lambda}}+n e^{\lambda(n-1)} \frac{1-e^{\lambda}}{1-e^{\lambda n}}\right]-2 \frac{1-e^{\lambda m}}{1-e^{\lambda}}\right\}^{-1} \tag{9}
\end{align*}
$$

Thus, the DTG unstable equilibrium state evaluation algorithm can be represented either by the system solution (8) or by calculation expression (9).

We will study the evaluation algorithm errors. To do this, we will find the solutions of differential equation (2) by the method of successive approximations when $\sin \alpha_{1} \approx \alpha_{1}-\frac{\alpha_{1}^{3}}{3!}$. At first, we introduce the replacement $x=\alpha_{1}$.

We write the equation

$$
\begin{equation*}
\varepsilon(t)=k_{\partial} t \tag{10}
\end{equation*}
$$

where $k_{\partial}$ - slope drift, taking into account the remarks in the first approximation :

$$
\begin{equation*}
\ddot{x}_{1}+2 \xi \dot{x}_{1}-\omega_{0}^{2} x=0 \tag{11}
\end{equation*}
$$

The solution of the equation of the first approximation has the form

$$
\begin{equation*}
x_{1}=\left(c_{10} e^{\omega_{10} t}+c_{2} e^{-\omega_{0} t}\right) e^{-\xi t} \tag{12}
\end{equation*}
$$

We write the equation of the second approximation

$$
\begin{equation*}
\ddot{x}_{2}+2 \xi \dot{x}_{2}-\omega_{0} x_{2}=-\frac{x_{1}^{3}}{6} . \tag{13}
\end{equation*}
$$

Taking into account the smallness of the second approximation in equation (13), set $\xi=0$; considering the expression (12) we find the solution of the second approximation from the equation
$\ddot{x}_{2}-x_{1} \omega_{0}^{2}=-\frac{1}{6}\left(c_{10}^{3} e^{3 \omega_{0} t}+3 c_{10}^{2} c_{20} e^{\omega_{0} t}+3 c_{10} c_{20}^{2} e^{-\omega_{0} t}+c_{20}^{3} e^{-3 \omega_{0} t}\right)$.
From equation (14) we obtain the second approximation solution
$x_{2}=-\frac{1}{48}\left(c_{10}^{3} e^{3 \omega_{0} t}+c_{20}^{3} e^{-3 \omega_{0} t}\right)-\frac{c_{10} c_{20} \omega_{0} t}{4}\left(c_{10} e^{\omega_{0} t}-c_{20} e^{\omega_{0} t}\right)$.
Taking into account the second approximation (expression (15) and Maclaurin series decomposition in the parameter $\xi_{1,}$

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find the law of SE motion nearby the state of unstable equilibrium
$\alpha(t)=R_{0}^{S}+c_{10} e^{\omega_{0} t}+c_{20} e^{-\omega_{0} t}-\xi t\left(c_{10} e^{\omega_{0} t}+c_{20} e^{-\omega_{0} t}\right)-$
$-\frac{1}{48}\left(c^{3}{ }_{10} e^{3 \omega_{0} t}+c^{3}{ }_{20} e^{-3 \omega_{0} t}\right)-\frac{c_{10} c_{20} \omega_{0} t}{4}\left(c_{10} e^{\omega_{0} t}-c_{20} e^{-\omega_{0} t}\right)+\varepsilon(t)$.
On the other hand, having expressed sh $\omega_{0} t$, $\operatorname{ch} \omega_{0} t$ through exponential functions of $e^{\omega_{0} t}, e^{-\omega_{0} t}$ type and having put the model functions (4) in Taylor series according to the parameters $\hat{R}^{S}, \hat{c}_{1}, \hat{c}_{2}, \omega_{0}$ in the neighborhood of the point $\hat{R}_{0}^{S}, \hat{C}_{10}, \hat{C}_{20}, \hat{\omega}_{0}$, keeping the first two expansion terms, we obtain
$\alpha(t)=R_{0}^{S}+\Delta R^{S}+\hat{c}_{10} e^{\omega_{0} t}+\hat{c}_{20} e^{-\omega_{0} t}+\Delta \hat{c}_{1} e^{\omega_{0} t}+\Delta \hat{c}_{2} e^{-\omega_{0} t}+$
$+\Delta \omega t\left(\hat{c}_{10} e^{\omega_{0} t}-\hat{c}_{20} e^{-\omega_{0} t}\right)$,
where $\Delta \hat{R}^{S}, \Delta \hat{c}_{1}, \Delta \hat{c}_{2}$ are state evaluation errors; $\Delta \hat{R}^{S}=\hat{R}^{S}-\hat{R}_{0}^{S}, \Delta \hat{c}_{1}=\hat{c}_{1}-c_{10}, \Delta \hat{c}_{2}=\hat{c}_{2}-c_{20} ; \Delta \omega=\omega-\omega_{0} \quad$ is the error of $\omega_{0}$.

Putting expressions (16) and (17) into the functional (5), we obtain
$F_{S}=\sum_{i=1}^{n}\left[\Delta \hat{R}^{s}+\Delta \hat{c}_{1} e^{\omega_{0} t_{i}}+\Delta \hat{c}_{2} e^{-\omega_{0} t_{i}}-\alpha_{S}^{0}\left(t_{i}\right)\right]^{2}$,
where
$\alpha_{S}^{0}\left(t_{i}\right)=-\Delta \omega t_{i}\left(c_{10} e^{\omega_{0} t_{i}}-c_{20} e^{-\omega_{0} t_{i}}\right)-\xi_{1} t_{i}\left(c_{10} e^{\omega_{0} t_{i}}+c_{20} e^{-\omega_{0} t_{i}}\right)-$
$-\frac{1}{48}\left(c_{10}^{3} e^{3 \omega_{0} t_{i}}+c_{20}^{3} e^{-3 \omega_{0} t_{i}}\right)-\frac{c_{10} c_{20} \omega_{0} t_{i}}{4}\left(c_{10} e^{\omega_{0} t_{i}}-c_{20} e^{-\omega_{0} t_{i}}\right)+\varepsilon\left(t_{i}\right)$.
Terms of achieving the functional minimum are equivalent to the system of equations

$$
\begin{equation*}
c_{\lambda_{0}, n}^{S} \hat{\delta}^{S}=\bar{f}^{S} \tag{20}
\end{equation*}
$$

where
$\hat{\delta}^{S}=\left[\Delta \hat{R}^{S} \Delta \hat{c}_{1} \Delta \hat{c}_{2}\right]^{T}, \quad \hat{f}^{S}=\left[f_{1}^{S} f_{2}^{S} f_{3}^{S}\right]^{T}=\left[\begin{array}{l}\sum_{i=1}^{n} \alpha_{S}^{0}\left(t_{i}\right) \sum_{i=1}^{n} \alpha_{S}^{0}\left(t_{i}\right) e^{\lambda_{0}(i-1)} \times \\ \times \sum_{i=1}^{n} \alpha_{S}^{0}\left(t_{i}\right) e^{-\lambda_{0}(i-1)}\end{array}\right]^{T}$,
$c_{\lambda_{0}, n}^{S}=\left.c_{\lambda, n}^{S}\right|_{\lambda=\lambda_{0}}, \quad \lambda_{0}=\omega_{0} \Delta t$.
Put the expression (9) into the system (20) and find the evaluation error

$$
\begin{align*}
& \Delta \hat{R}^{S}=\left\{f_{1}^{S}\left[\frac{1+e^{\lambda_{0} u}}{1+e^{\lambda_{0}}}+n e^{\lambda_{0}(n-1)} \frac{1-e^{\lambda_{0}}}{1-e^{\lambda_{0} n}}\right]-f_{2}^{S}-f_{3}^{S} e^{\lambda_{0}(n-1)}\right\} \times  \tag{21}\\
& \times\left\{n\left[\frac{1+e^{\lambda_{0} n}}{1+e^{\lambda_{0}}}+n e^{\lambda_{0}(n-1)} \frac{1-e^{\lambda_{0}}}{1-e^{\lambda_{0} n}}\right]-2 \frac{1-e^{\lambda_{0} n}}{1-e^{\lambda_{0}}}\right\} .
\end{align*}
$$

Having made the boundary transition at $\Delta t \rightarrow 0$ and $T_{H}=$ const, we obtain an expression of the evaluation error

$$
\begin{equation*}
\Delta \hat{R}^{S}=\sum_{j=1}^{S} \Delta R_{j}^{S} \tag{22}
\end{equation*}
$$

where

$$
\begin{align*}
& \Delta \hat{R}^{S}=-\frac{1}{48}\left(c_{10}^{3} e^{\frac{3 u}{2}}+c_{20}^{3} e^{-\frac{3 u}{2}}\right) \cdot l_{2}(u),  \tag{23}\\
& \Delta \hat{R}_{2}^{S}=-\frac{\xi_{1}}{\omega_{0}} l_{1}(u)\left(c_{10} e^{\frac{u}{2}}-c_{20} e^{-\frac{u}{2}}\right),  \tag{24}\\
& \Delta \hat{R}_{3}^{S}=-\frac{c_{10} c_{20}}{4}\left(c_{10} e^{\frac{u}{2}}+c_{20} e^{-\frac{u}{2}}\right) \cdot l_{1}(u),  \tag{25}\\
& \Delta \hat{R}_{4}^{S}=-\frac{\Delta \omega}{\omega_{0}} l_{1}(u)\left(c_{10} e^{\frac{u}{2}}+c_{20} e^{-\frac{u}{2}}\right),  \tag{26}\\
& \Delta \hat{R}_{5}^{S}=\omega_{0} \Delta^{-1} \int_{0}^{T}\left\{\left[\frac{1}{2}\left(e^{u}+1\right)+\frac{u e^{u}}{e^{u}-1}\right]-e^{\omega_{0} t}-e^{u} e^{-\omega_{0} t}\right\} \varepsilon(t) d t,  \tag{27}\\
& l_{1}(u)=e^{\frac{u}{2}} \Delta^{-1}\left[\frac{u e^{u}}{2}\left(\frac{u\left(e^{u}+1\right)}{e^{u}-1}-1\right)-\frac{e^{2 u}-1}{4}\right],  \tag{28}\\
& l_{2}(u)=\frac{1}{3} e^{-\frac{3 u}{2}} \Delta^{-1}\left[u e^{u}\left(e^{2 u}+e^{u}+1\right)-\frac{1}{4}\left(e^{2 u}-1\right)\left(e^{2 u}+4 e^{u}+1\right)\right],  \tag{29}\\
& \Delta=u\left(\frac{e^{u}+1}{2}+\frac{u e^{u}}{e^{u}-1}\right)-2\left(e^{u}-1\right), u=\omega_{0} T_{c},
\end{align*}
$$

$\Delta \hat{R}_{1}^{S}$ is the component of the evaluation error caused in the law of motion of SE by the item $-\frac{1}{48}\left(c_{10}^{3} e^{3 \omega_{0} t}+c_{20}^{3} e^{-3 \omega_{0} t}\right)$; $\Delta \hat{R}_{2}^{S}$ is the component of the evaluation error caused by neglecting the decrement of attenuation in the model (4); $\Delta \hat{R}_{3}^{S}$ is the component of the evaluation error caused in the law of motion of SE by the item $-\frac{1}{4} c_{10} c_{20} \omega_{10} t\left(c_{10} e^{\omega_{0} t}-c_{20} e^{-\omega_{0} t}\right)$; $\Delta \hat{R}_{4}^{S}$ is the component of the evaluation error caused by the error of setting the parameter $\omega_{0} ; \Delta \hat{R}_{5}^{s}$ is the component of the evaluation error caused by the distortions of the observed law of motion of SE .

We will analyze the components of the evaluation error $\Delta \hat{R}_{j}^{S}(j=1,5)$.

As we can see from the obtained analytical expressions of these components of the error, their values depend on the parameters $\frac{\Delta \omega}{\omega_{0}}, \frac{\xi_{1}}{\omega_{0}}$, the initial conditions $\alpha_{0}, \dot{\alpha}_{0}$, and the information surveillance time. To analyze the maximums of the components $\Delta \hat{R}_{j}^{S}$ we define a zone of initial conditions on the phase plane , located inside the circle of radius A, that is

$$
\begin{equation*}
\alpha_{0}^{2}+\left(\frac{\dot{\alpha}_{0}}{\omega}\right)^{2} \leq A^{2} \tag{30}
\end{equation*}
$$

We use the polar coordinate system on the phase plane $O \alpha \dot{\alpha}$ (Fig. 1). We combine plus and polar axis with $O$ start and abscissa axis $O \alpha$ of the rectangular coordinate system.

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We denote $\rho$ and $\varphi$ as polar radius and polar angle respective$l y$, and the equation is

$$
\begin{equation*}
\alpha_{0}=\rho \cos \varphi, \frac{\dot{\alpha}}{\omega}=\rho \sin \varphi \tag{31}
\end{equation*}
$$

Then the constants of integration $C_{10}, C_{20}$ are expressed with the use of $\rho$ and $\varphi$ by such relations

$$
\begin{equation*}
c_{10}=\frac{\rho}{\sqrt{2}} \sin \left(\varphi+\frac{\pi}{4}\right), c_{20}=\frac{\rho}{\sqrt{2}} \cos \left(\varphi+\frac{\pi}{4}\right) \tag{32}
\end{equation*}
$$

Fig. 1. Polar coordinate system
Taking into account (32), the expressions (23) ... (26) are written as .

$$
\begin{align*}
& \Delta R_{1}^{S}=-\frac{\rho^{3}}{96 \sqrt{2}} l_{2}(u)\left[e^{\frac{3 u}{2}} \sin ^{3}\left(\varphi+\frac{\pi}{4}\right)+e^{-\frac{3 u}{2}} \cos ^{3}\left(\varphi+\frac{\pi}{4}\right)\right],  \tag{33}\\
& \Delta \hat{R}_{2}^{S}=-\frac{\xi_{1}}{\omega_{0}} \rho l_{1}(u) \sqrt{\operatorname{ch} u} \sin \left(\varphi+\frac{\pi}{4}-\operatorname{arctg} e^{-u}\right),  \tag{34}\\
& \Delta \hat{R}_{3}^{S}=-\frac{\rho}{16} l_{1}(u) \sqrt{\operatorname{ch} u} \sin 2\left(\varphi+\frac{\pi}{4}\right) \sin \left(\varphi+\frac{\pi}{4}+\operatorname{arctg} e^{-u}\right),  \tag{35}\\
& \Delta \hat{R}_{4}^{S}=-\frac{\Delta \omega}{\omega} \rho l_{1}(u) \sqrt{\operatorname{chu}} \sin \left(\varphi+\frac{\pi}{4}+\operatorname{arctge}^{-u}\right) . \tag{36}
\end{align*}
$$

Based on the components of expressions (33) ... (36), we obtain the maximum values of evaluation error if $\rho=A$ :

$$
\begin{aligned}
& \max \left(\Delta \hat{R}_{1}^{s}\right)=\frac{A^{3}}{192} l_{2}(u) e^{\frac{3 u}{2}} \\
& \max \left(\Delta \hat{R}_{2}^{S}\right)=-\frac{\left|\xi_{0}\right|}{\omega_{0}} A l_{1}(u) \sqrt{c h u} \\
& \max \left(\Delta \hat{R}_{3}^{S}\right)=-\frac{A^{3}}{16} l_{1}(u) \sqrt{c h u} \cdot m(u) \\
& \max \left(\Delta \hat{R}_{4}^{S}\right)=-\frac{\Delta \omega}{\omega_{0}} A l_{1}(u) \sqrt{c h u}
\end{aligned}
$$

where $m(u)=\left|\sin 2\left(\varphi+\frac{\pi}{4}\right) \sin \left(\varphi+\frac{\pi}{4}+\operatorname{arctg} e^{-u}\right)\right|$,

$$
\begin{aligned}
& \varphi=-\frac{\pi}{4}-\operatorname{arctg}\left[2 \sqrt{\frac{p}{3}} \cos \frac{y}{3}\right], \cos y=\frac{q}{2 \sqrt{\left(\frac{p}{3}\right)^{3}}} \\
& p=2+\frac{4}{3} e^{-2 u}, q=\frac{e^{-u}}{3}+\frac{16}{27} e^{-3 u}
\end{aligned}
$$

We select from the group of error components expressions (37) the coefficients that determine the nature of the changes of the maximum values of these components:

$$
\begin{align*}
& v_{1}(u)=l_{2}(u) e^{\frac{3 u}{2}} \\
& v_{2}(u)=-l_{1}(u) \sqrt{c h u}=v_{4}(u),  \tag{38}\\
& v_{3}(u)=-l_{1}(u) \sqrt{c h u} \cdot m(u)=v_{2}(u) \cdot m(u)
\end{align*}
$$

Dependence of the coefficients (38) on information surveillance time when changing $T_{H}$ from 0 to $T_{0}$ with steps $0,005 T_{0}$ is given in the form of graphs (Fig. 2). As follows from these relationships, the coefficients increase with the increase of information surveillance time.


Fig. 2. Dependence of the impact of errors $v_{1}(\mathrm{a}), v_{2} v_{3}(\mathrm{~b})$ on $T_{c}$ $\left(1-v_{2} ; 2-v_{3}\right)$

According to these dependences of the coefficients, dependences of maximum values of error components (Fig. 3 a , b) have been made. Based on these charts, the following conclusions can be made:

- for the zone of the initial conditions (30) the maximum error value $\Delta \hat{R}_{1}^{S}$ increases when information surveillance time increases and is determined by the parameter $A$. So, if $A=1^{\circ}$ and $T_{c}=0,3 T_{0} \max \left(\Delta \hat{R}_{1}^{s}\right)=1,33$ Arcsecond (Fig. 3, a). According to (37), based on a cubic dependence $\max \left(\Delta \hat{R}_{1}^{S}\right)$ on $A$, by decreasing the radius of the zone of initial conditions arbitrarily small impact of this component on the total evaluation error at the information surveillance time that does not exceed 0,3 $T_{0}$ can be achieved;



Fig. 3. Dependence of errors $\left|\Delta R_{i}^{S}\right|$ on $T_{c}$ where i: 1-1; 2-3; 3-2; 4

- the maximum values of evaluation errors $\Delta \hat{R}_{2}^{s}$ and $\Delta \hat{R}_{4}^{S}$ are going up with increasing of information surveillance time. Taking into account that in real devices $(\Delta \omega) \omega_{0}^{-1} \leq 5 \cdot 10^{-4}$ and $\left|\xi_{1}\right| \cdot \omega_{0}^{-1}=10^{-4}$ when $A=1^{\circ}$ and $T=0,3 T_{0}$ we obtain that $\max \left(\Delta \hat{R}_{2}^{S}\right)=1,2 \operatorname{Arcsecond}$ and $\max \left(\Delta \hat{R}_{4}^{s}\right)=7$ Arcsecond. Based on the linear dependence of errors on $\rho$, the errors can be reduced to any small values by choosing $\rho$;
- if $\rho \leq 1^{\circ}$ for information surveillance time $T_{H}=0,3 T_{0}$, the impact of the component $\Delta \hat{R}_{3}^{S}$ on the total error $\Delta \hat{R}^{s}$ is negligibly small (see Fig. 3).
On the other hand, according to (33) ... (36), in the case of preassigned $T_{c}$ for each evaluation error component there are the optimal polar angles, in the sense of transformation of this component into zero.

For evaluation error $\Delta \hat{R}_{1}^{s}=0$ the optimal polar angle is

$$
\begin{equation*}
\varphi_{\text {AnT }}=k \pi-\frac{\pi}{4} \tag{39}
\end{equation*}
$$

$\mathrm{k}=0 ; \pm 1 ; \pm 2 \ldots$.
Evaluation error component $\Delta \hat{R}_{2}^{S}$ becomes zero at

$$
\begin{equation*}
\varphi_{s n T}=k \pi-\frac{\pi}{4}+\operatorname{arctg} e^{-u} \tag{40}
\end{equation*}
$$

$k=0 ; \pm 1 ; \pm 2 \ldots$
Evaluation error component $\Delta \hat{R}_{3}^{S}=0$ at

$$
\begin{equation*}
\varphi_{M n T}=k \pi-\frac{\pi}{4} \text { i } \varphi_{s n T}=k \pi-\frac{\pi}{4}-\operatorname{arctg} e^{-u} \tag{41}
\end{equation*}
$$

$k=0 ; \pm 1 ; \pm 2 \ldots$.
Evaluation error component becomes zero at

$$
\varphi_{м n T}=k \pi-\frac{\pi}{4}-\operatorname{arctg} e^{-u}
$$

$k=0 ; \pm 1 ; \pm 2 \ldots$.
In (39) ... (42) the following denotations are accepted

$$
\begin{equation*}
u=\omega_{0} T_{H}, \operatorname{tg} \varphi=\frac{\dot{\alpha}_{0}}{\omega_{0} \alpha_{0}} \tag{43}
\end{equation*}
$$

Consider the impact of the typical distortions on the accuracy of evaluation of DTG state for the observed law of motion of SE .

1. Linear drift of SE equilibrium, as a result of harmful moments in the case of southern orientation of SE, leads to the error of stable equilibrium (Fig. 4).


Figure 4. Dependence of $\Delta \hat{R}_{52}^{S}$ on such distortions as $e^{-\frac{t}{T}}$ and $T_{c}$ $\left(\mathrm{a}_{\mathrm{e}}=1{ }^{\prime}\right)$ at $\tau, \mathrm{s}: 1-300 ; 2-100$
2. As it is proved in [9], SE is affected by the moments with a period comparable or much smaller than the period of precession oscillations. In both cases the distortion is given by the expression
$\Delta \hat{R}_{53}^{N}=a_{r} d_{33}(\mu, u) \sin \left(\frac{\mu u}{2}+\varphi_{1}\right)$,
After substituting (44) in the expression of evaluation error $\Delta \hat{R}_{5}^{S}$ and limiting transition at $\Delta \mathrm{t} \rightarrow 0, \quad T_{c}=$ const, we can write an expression of evaluation error

$$
\begin{equation*}
\Delta \hat{R}_{53}^{S}=a_{r} l_{3}(\mu, u) \sin \left(\frac{\mu u}{2}+\varphi_{1}\right) \tag{45}
\end{equation*}
$$

where
$I_{3}(\mu, u)=\frac{2}{\mu} \cdot \frac{\left(\frac{e^{u}+1}{2}+\frac{u e^{u}}{e^{u}-1}\right) \sin \frac{\mu u}{2}-\frac{2}{\mu^{2}+1}\left[\left(e^{u}-1\right) \cos \frac{\mu u}{2}+\mu\left(e^{u}+1\right) \sin \frac{\mu u}{2}\right]}{u\left(\frac{e^{u}+1}{2}+\frac{u e^{u}}{e^{u}-1}\right)-2\left(e^{u}-1\right)}$.
Graphical dependency of the coefficient $l_{3}(\mu, u)$ for $\mu=2,3$, 5 is shown in Fig. 5. We can see that $l_{3}(\mu, u)$ reduces from a certain value, which depends on $\mu$, to zero, and then repeatedly crosses the x-axis. Graphical dependency of errors (Fig. 6) shows that the state evaluation error $\Delta \hat{R}_{53}^{S}$, as well as the coefficient $l_{3}(\mu, u)$, varies according to the harmonic law. Lowfrequency disturbance at $T_{c} \leq 0,3 T_{0}$ is weakly subdued by the state evaluation algorithm. The increase in information surveillance time $T_{c}>0,15 T_{0}$ and interference frequency reduces the evaluation error. If $\mu=100$ and $T_{c} \leq 0,1 T_{0}$, the interference amplitude is hundreds of times smaller than the evaluation error; thus there has been increasing interference that under these conditions is greater when the noise frequency is higher. With increasing $T_{c} \geq 0,15 T_{0}$ the error amplitude is subdued in hundreds of times (Fig. 6, a) and at $T_{c} \geq 0,2-0,3 T_{0}$ the high frequency harmonic distortion $\Omega \geq 100 \omega 0$ hardly causes the evaluation error.


Fig. 5. Dependence of the impact coefficient $I_{3}$ on $T_{c}$, when $\mu$ : 1 - 100; 2-5; 3-3; 4-2

Thus, the low-frequency distortions in the SE motion law at $0,1 \mathrm{~T}_{0} \leq \mathrm{T}_{\mathrm{c}} \leq \mathrm{T}_{0}$ are weakly subdued by the state evaluation
algorithm, and high-frequency distortions $\left(\Omega=100 \omega_{0}\right)$ are effectively filtered at $T_{c} \geq 0,15 T_{0}$, and at $T_{c}=0,2 \ldots 0,3 T_{0}$ they barely cause the evaluation error.


Fig. 6. Dependence of error $\Delta \hat{R}_{53}^{S}$ on distortions such as $\sin \Omega t, T_{c}$ $\left(a_{r}=1\right.$ '), when $\mu$ : $a-10$ and 5 (respectively curves 1 and 2); b-3; 2 and 0.5 ( respectively curves 1,2 and 3 )
3. If the analysis of evaluation error because of the noise meter (AS) in the observed law of motion of SE as well as in [9], we assume that measuring error is given in the form of white noise and for evaluation error we use the Holder inequality. Under these conditions evaluation can be represented by the expression

$$
\begin{equation*}
\left|\Delta \hat{R}_{54}^{S}\right| \leq l_{34}(u) \Phi_{\varepsilon}^{S}\left(T_{H}\right) \tag{46}
\end{equation*}
$$

where

$$
l_{34}(u)=\sqrt{\frac{u(u+s h u)}{u(u+s h u)-8 \operatorname{sh}^{2} u \cdot 2^{-1}}} ; \Phi_{\varepsilon}^{S}\left(T_{c}\right)=\sqrt{T_{H}^{-1} \int_{0}^{T_{c}} \varepsilon^{2}(t) d t} ;
$$

$\Phi_{\varepsilon}^{S}\left(T_{c}\right)$ is quadratic mean value of the random function on the interval $[0, \mathrm{~T}]$.

Just as in the previous case, the coefficient $l_{34}(u)$ decreases with increasing of $T_{c}$ and actually does not change at $T_{c} \geq 0,2-$ $0,3 T_{0}$, and further increase in information surveillance time is unnecessary (Fig. 7).


Fig. 7. Dependence of the impact coefficient $I_{34}$ on $T_{c}$

## Conclusions:

As the result of the research done in this paper, the state evaluation algorithm of unstable equilibrium of DTG was developed. Evaluation errors were examined.

The analysis of mathematical expressions of evaluation errors revealed that the evaluation error contains five components, based on:

- non-linear motion trajectory distortions of the SE in large deviations from the North-South flat;
- non-linear frequency deviations of the precession oscillations, which occur due to anisochronism of movements of the SE near the unstable equilibrium;
- the SE motion trajectory bending due to the moments such as viscous friction;
- the precession oscillations cyclic frequency divergency that is used in the state evaluation algorithm, and in DTG precession oscillations frequency algorithm;
- presence of the disturbance such as linear and exponential drift of equilibrium state of the SE, harmonic interferences and white noise interferences in the observable motion law.

The first two evaluation errors components are directly proportional to the cube of polar radius $\rho=A=\sqrt{\alpha_{0}^{2}+\left(\dot{\alpha}_{0} \omega^{-1}\right)^{2}}$, where $\alpha_{0}, \dot{\alpha}_{0}$ is the initial angle and angular velocity of the SE movement in the respect of the unstable equilibrium. Another two components are proportional to the polar radius. Optimal relations (in terms of evaluation error vanishing) between the information surveillance time and initial conditions $\alpha_{0}=$ const, $\dot{\alpha}_{0} \rightarrow 0$ were found. It is demonstrated that all of the enumerated components of the evaluation error grow with the increase of the
information surveillance time, when initial conditions are close to $\alpha_{0}=$ const and $\dot{\alpha}_{0} \rightarrow 0$.

The examinations of the disturbance influence on the evaluation precision show that:

- disturbance such as linear drift of equilibrium state of the SE causes errors proportional to the drift speed and information surveillance time;
- exponential drift disturbance causes the error, which is mon-otone-decreasing, when information surveillance time increases. The smaller the time constant of the exponent is, the faster the error decrease is;
- low-frequency harmonic interference, the frequency of which is commensurable with the precession oscillations frequency, is weakly eliminated by the state evaluation algorithm;
- high-frequency harmonic interference, the frequency of which 100 times exceeds the precession oscillations frequency, and random white noise interference are effectively filtered by the state evaluation algorithm. Information surveillance time needed for complete HF-interferences and white noise interferences elimination has to be $0.2 \ldots 0.3$ of the precession oscillation period.


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